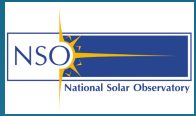


# Elementary Fluxtubes and Coronal Heating

Han Uitenbroek  
National Solar Observatory/Sacramento Peak

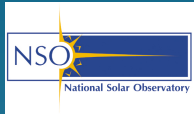


High-Resolution Magnetography from Space: Beyond Solar-B,  
Hunstville AL, Apr 3-5, 2001



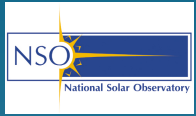
# Elementary Fluxtubes and Coronal Heating

- Introduction
- Magnetostatic model of network element
  - Hydrogen ionization
  - Partial frequency Redistribution (PRD) in Lyman- $\alpha$  is important
- Conclusions



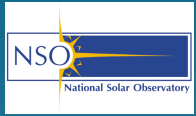
## Conclusion

# Multi-wavelength Stokes Spectroscopy and Sophisticated Radiative Transfer Modeling



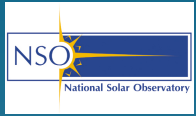
# Transition from Photosphere to Corona

- Difficult to observe
  - Highly structured in three-dimensions
  - Transparent in almost all diagnostics, in particular those accessible from the ground



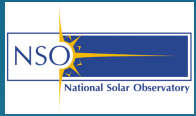
# Transition from Photosphere to Corona

- Difficult to observe
  - Highly structured in three-dimensions
  - Transparent in almost all diagnostics, in particular those accessible from the ground
- Difficult to model
  - Not stationary, but highly dynamic
  - Transition from gas pressure ( $\beta > 1$ ) to magnetic field ( $\beta < 1$ ) dominance
  - Transition from optically thick ( $\tau > 1$ ) to optically thin ( $\tau < 1$ )



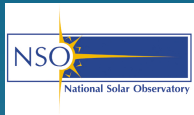
## To solve the problem of Coronal heating:

- How the structure of the small-scale solar magnetic field changes when we progress from the photosphere through the chromosphere up to the corona



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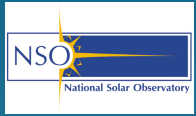
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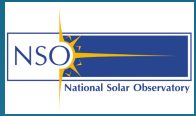
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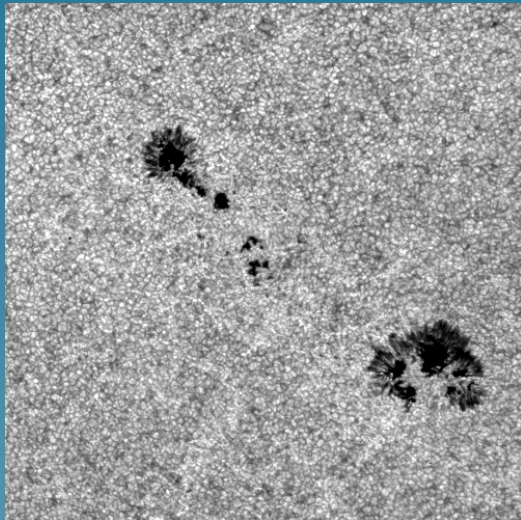


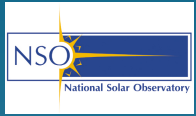
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- Is the field braided? Is energy transported via waves, and if so what kind of waves? Or is magnetic energy stored in the coronal field by continuous footpoint motions and released by reconnection?
- Can we explain the vertical persistence of the magnetic network?

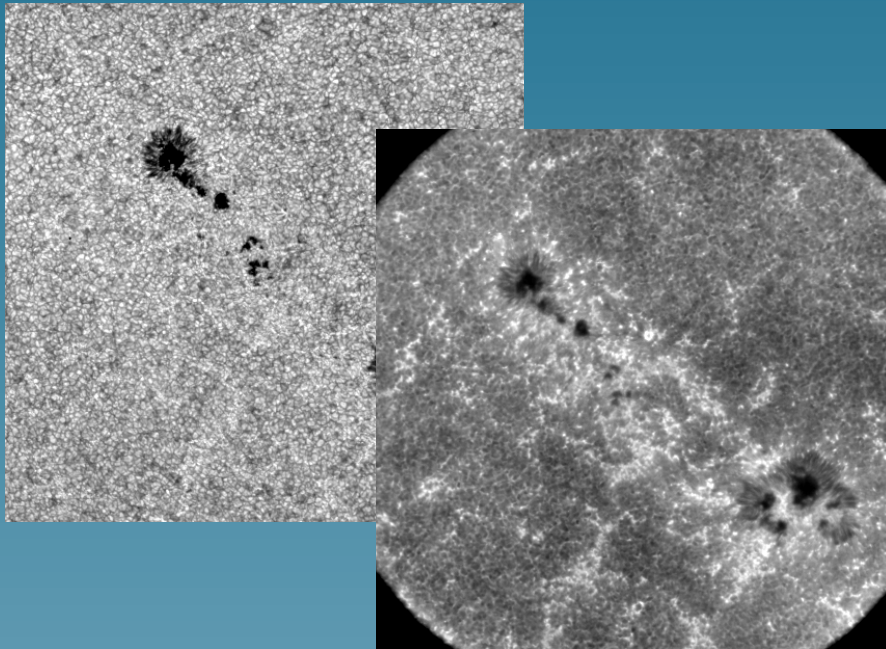


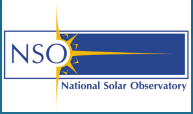
# Photosphere →



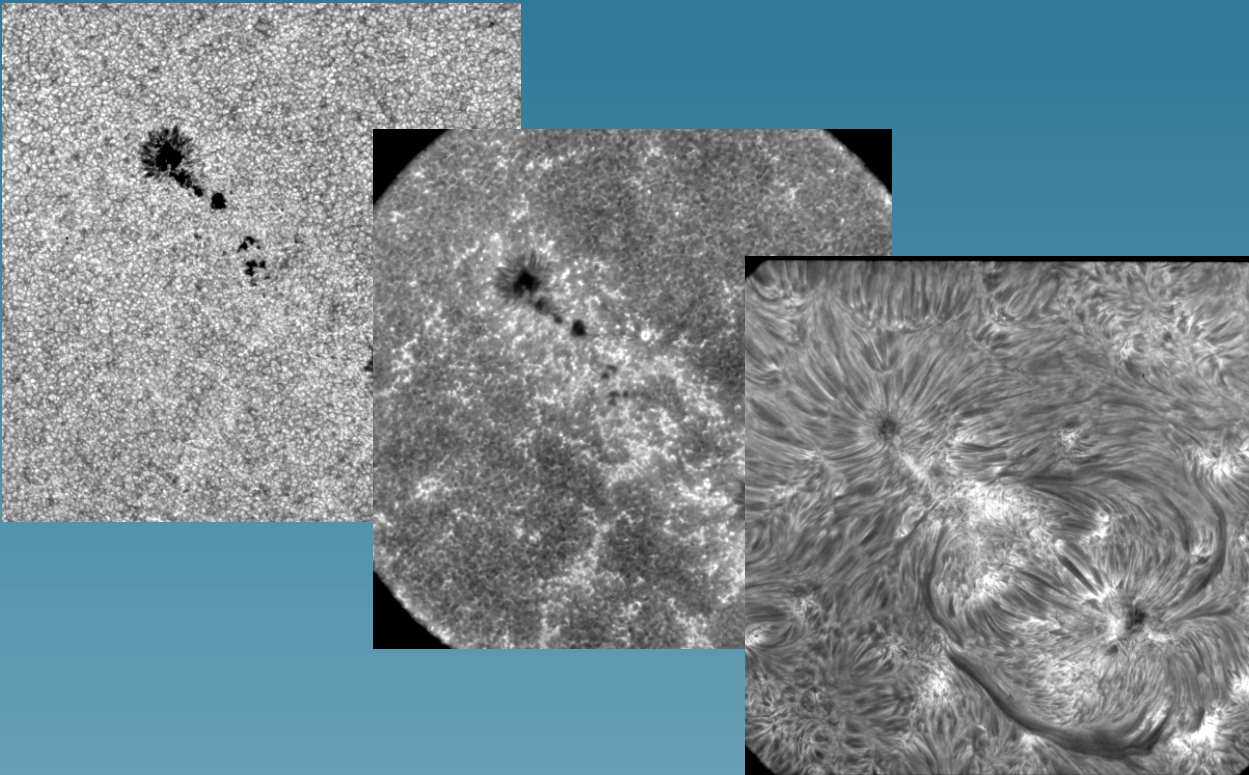


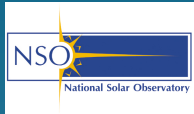
Photosphere → lower Chromosphere →



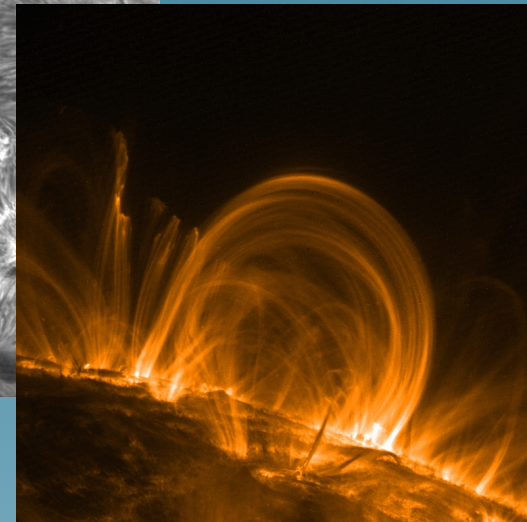
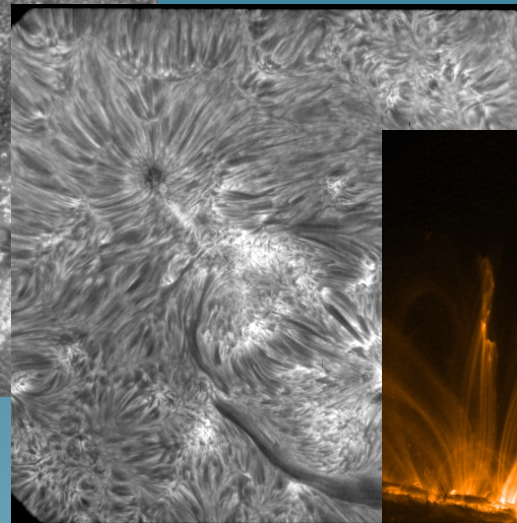
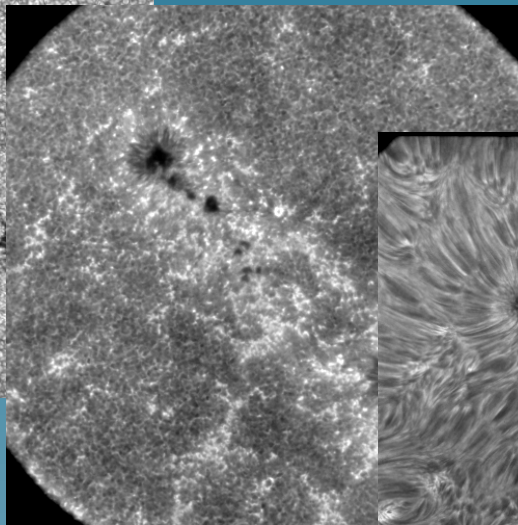
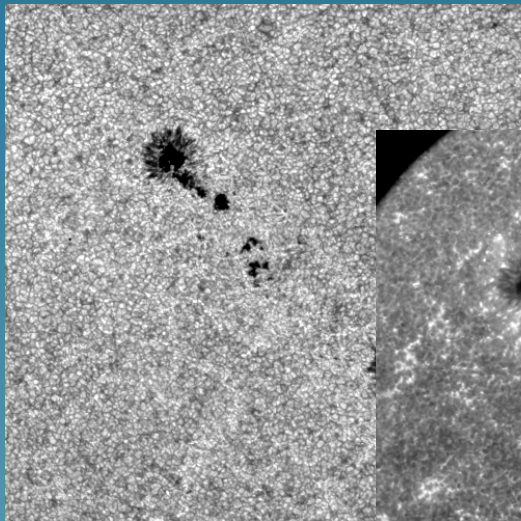


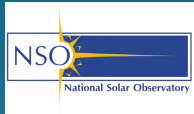
Photosphere → lower Chromosphere →  
upper Chromosphere →



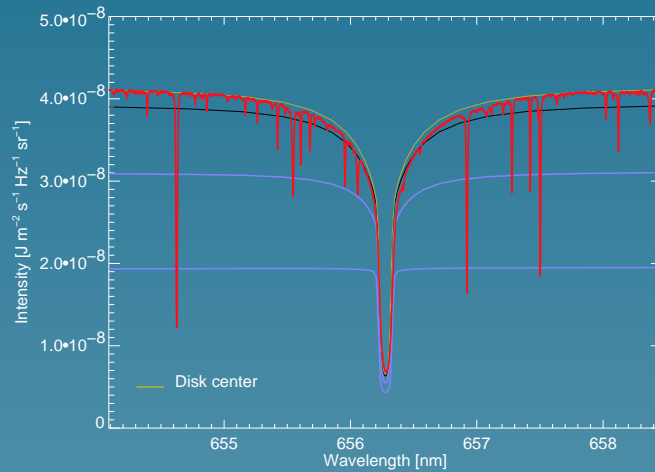


Photosphere → lower Chromosphere →  
upper Chromosphere → Corona



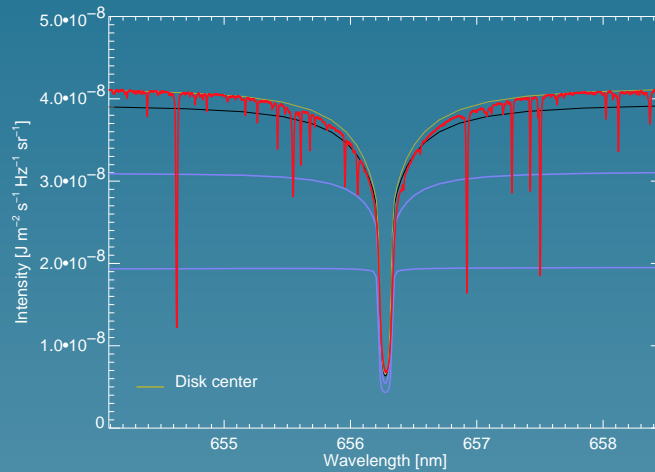


# H- $\alpha$ is important diagnostic

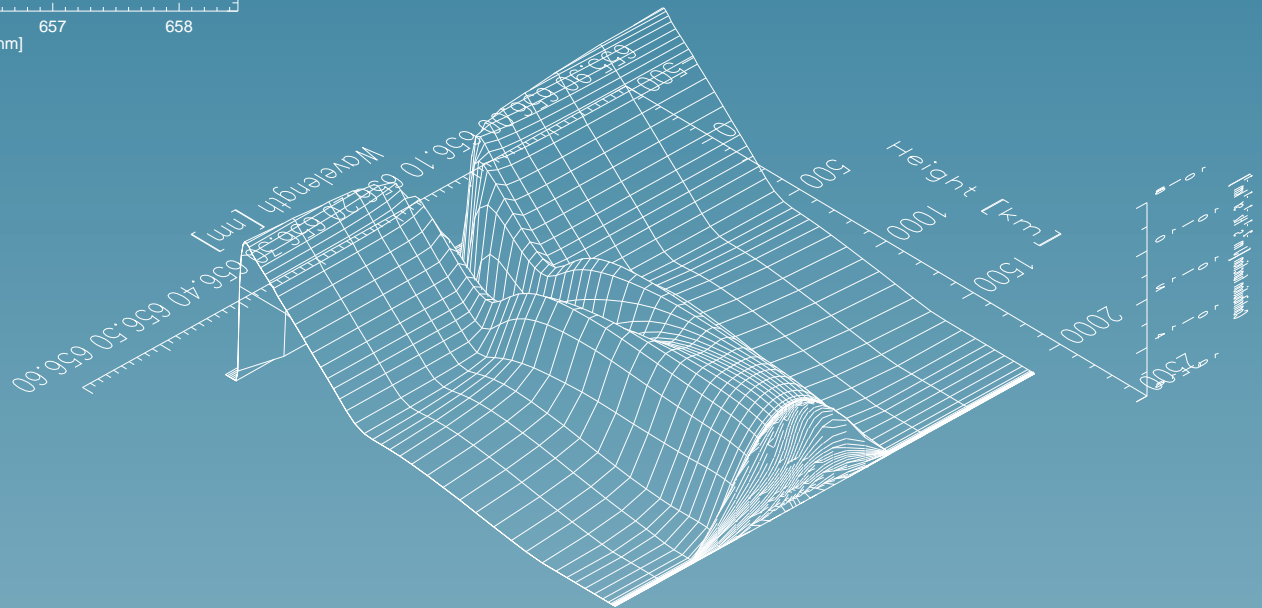




# H- $\alpha$ is important diagnostic

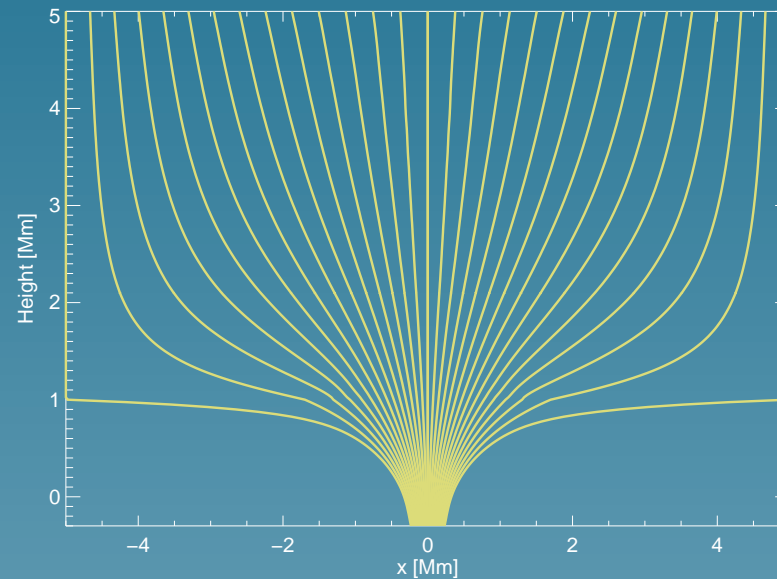


- Double-peaked **contribution function**
- Large velocities and steep flanks
- Scattering line core



## Two-dimensional network model

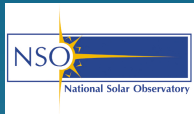
Magnetostatic models together with Aad van Ballegooijen (CfA)



Fluxtube in magnetostatic equilibrium:

$$-\nabla p + \rho g + (4\pi)^{-1}(\nabla \times B) \times B = 0$$





## Two-dimensional network model (2)

Pressure balance with the surroundings below the canopy:

$$p_{\text{int}}(z) + B_{\text{int}}^2(z)/8\pi = p_{\text{ext}}(z)$$

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Solution by minimization of the Lagrangian:

$$W = \int_{z_b}^{z_t} \int_0^{R(z)} \left( \frac{B^2}{8\pi} - p(\Phi, z) \right) 2\pi r dr dz + \int_{z_b}^{z_t} p_{\text{ext}}(z) \pi R^2(z) dz$$

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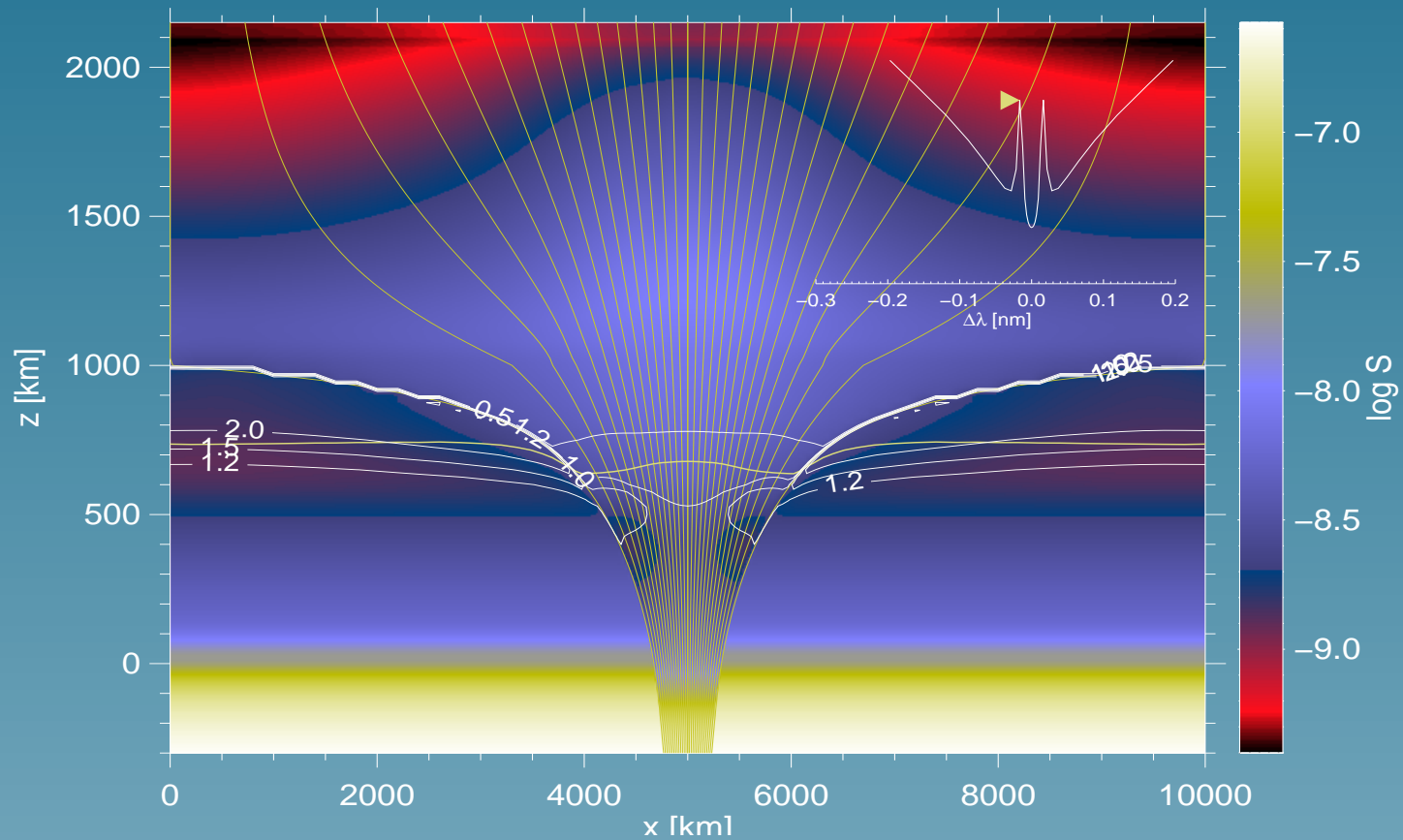
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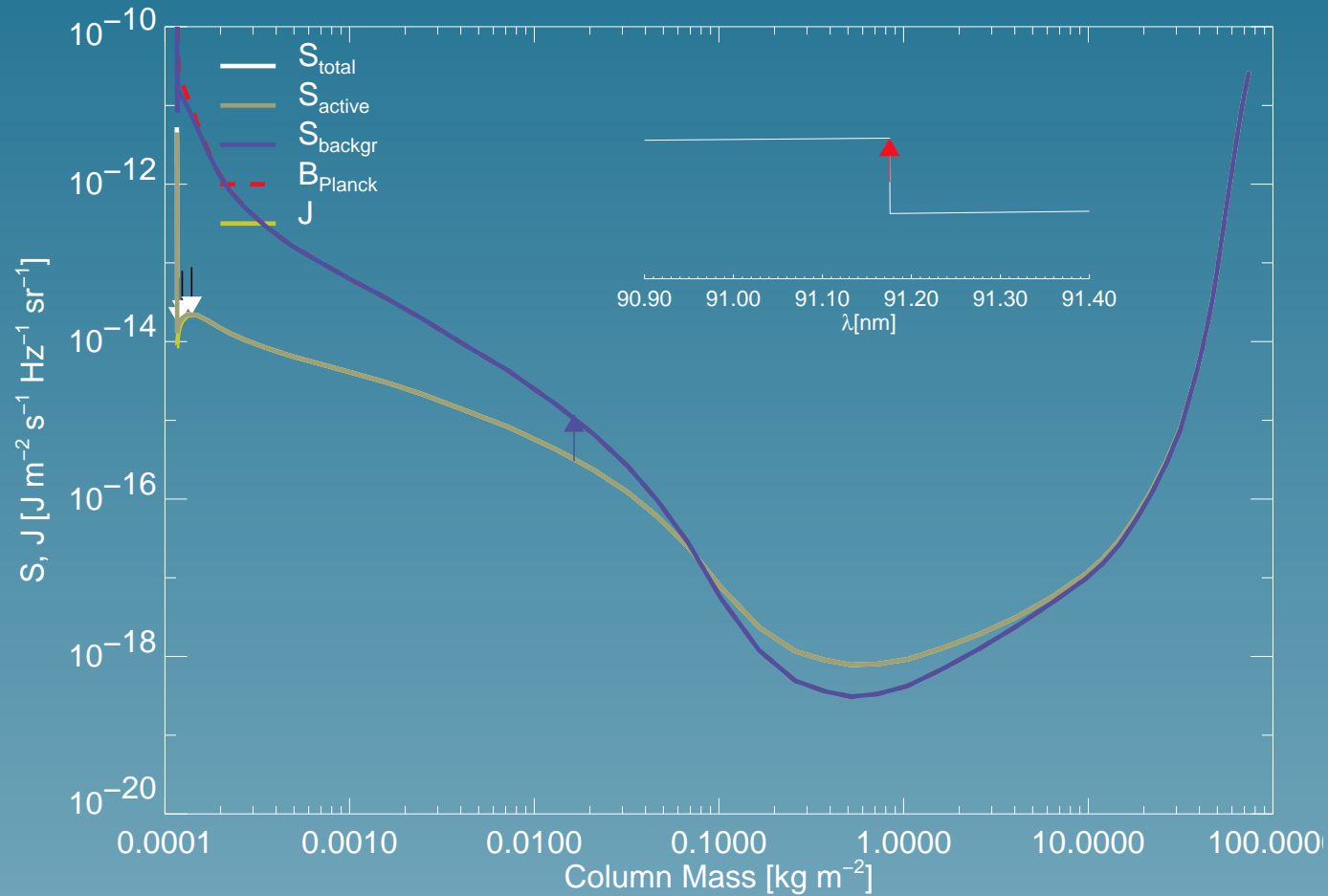
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Temperature  $T(r, z)$  and electron density  $N_e(r, z)$  for given  $p(r, z)$  and  $\rho(r, z)$  are computed consistent with global radiative transfer in hydrogen (Lyman and Balmer continua, Ly- $\alpha$ , Ly- $\beta$ , and H- $\alpha$ ).

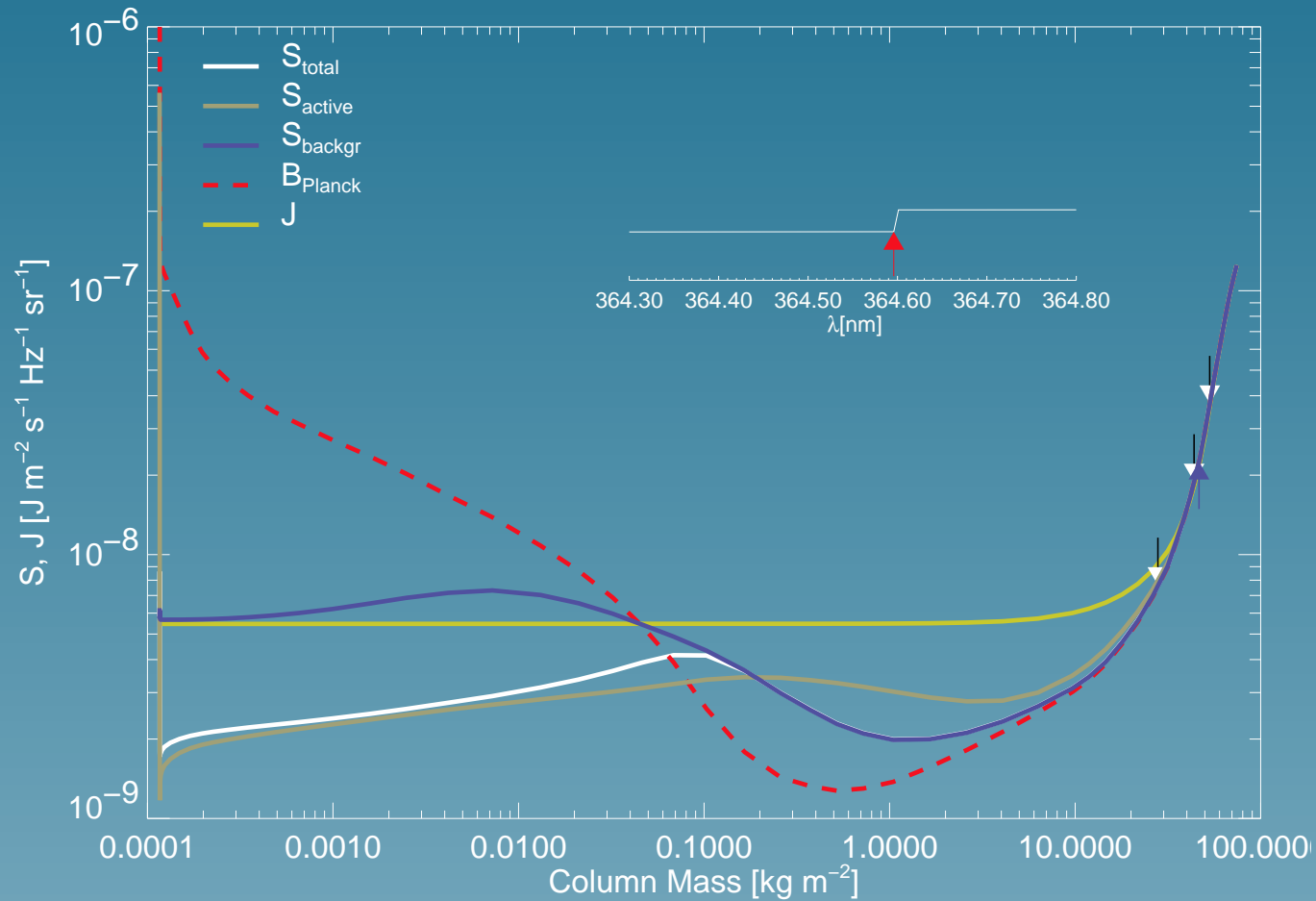
# Network Model



# Hydrogen ionization — Lyman continuum



# Hydrogen ionization — Balmer continuum

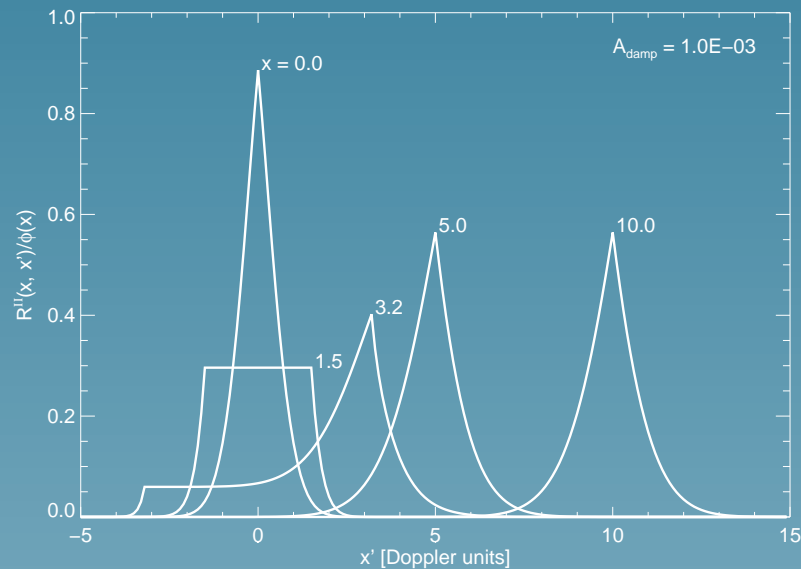


## Partial Frequency ReDistribution (PRD)

$$\frac{\psi_{ij}^{\text{PRD}}(\nu)}{\phi_{ij}(\nu)} = 1 + \frac{n_i B_{ij}}{n_j P_j} \int d\nu' J(\nu') \left[ \frac{R_{kji}(\nu, \nu')}{\phi_{ij}(\nu)} - \phi_{kj}(\nu') \right] \quad (1)$$

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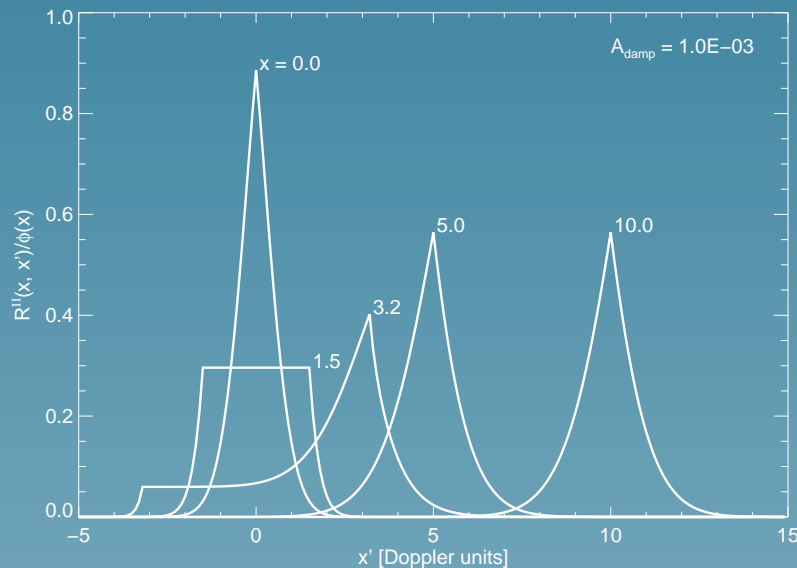
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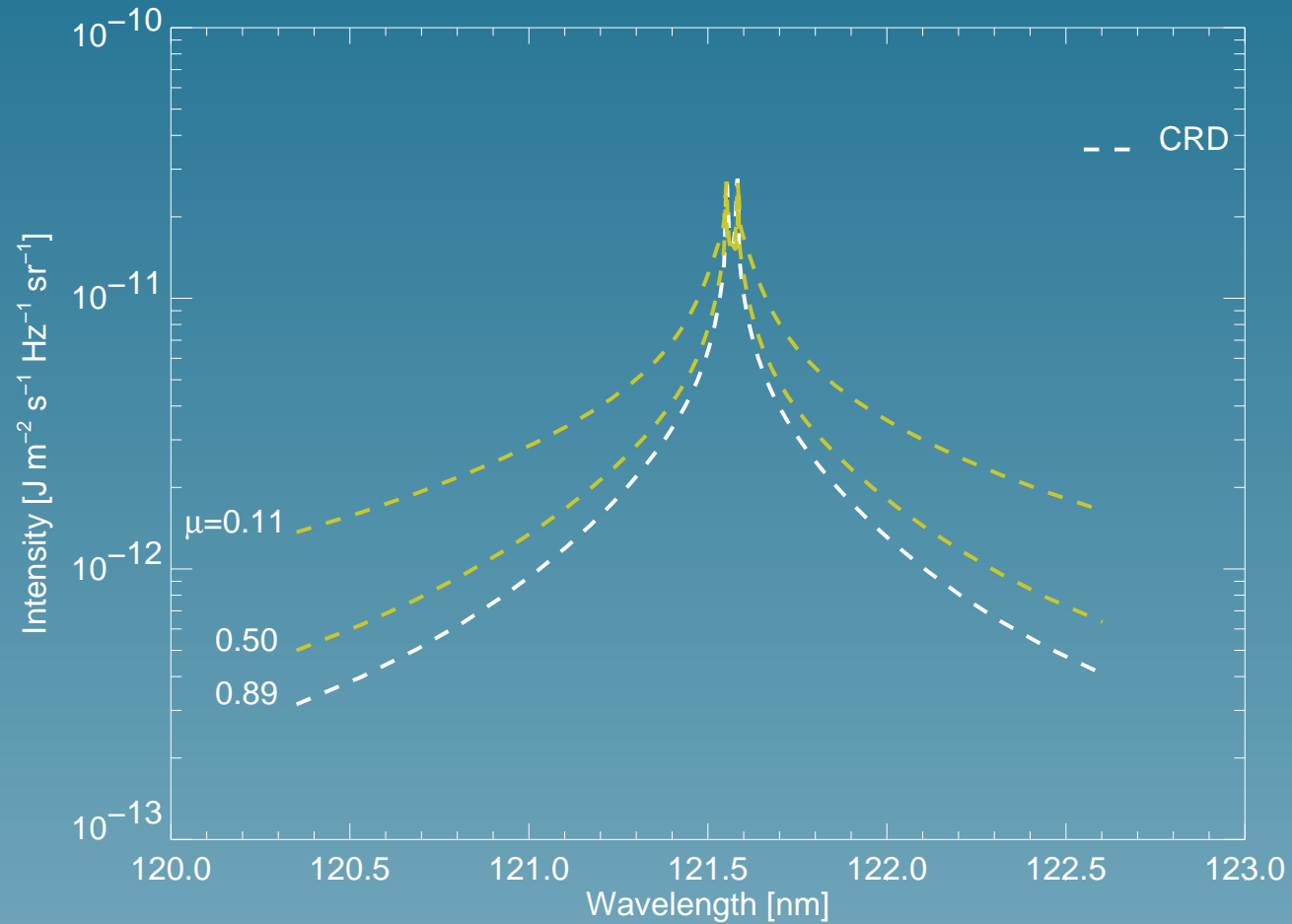
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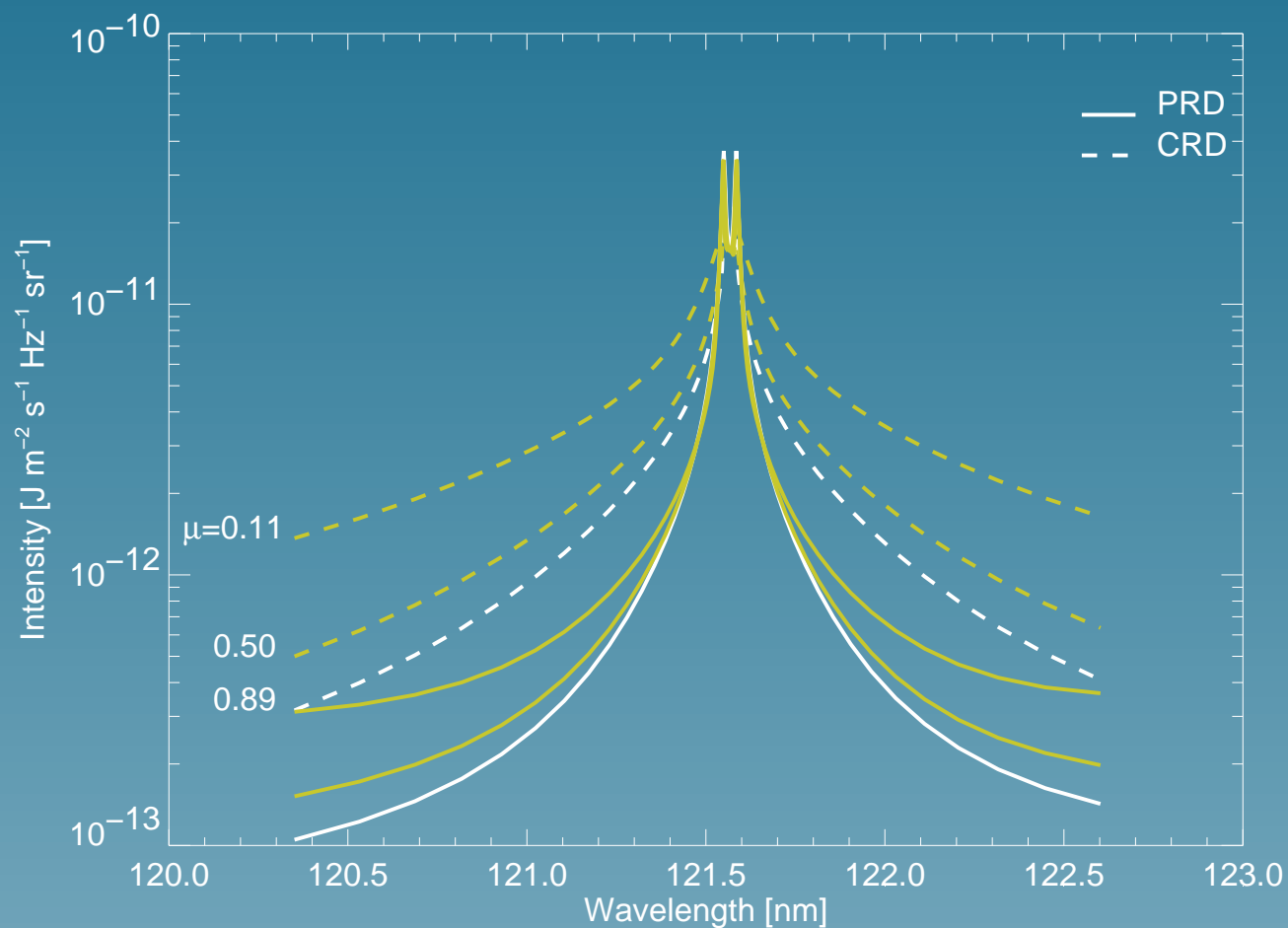


- Complete redistribution in core
- Coherent scattering in the wings
- Decoupling of wing source function

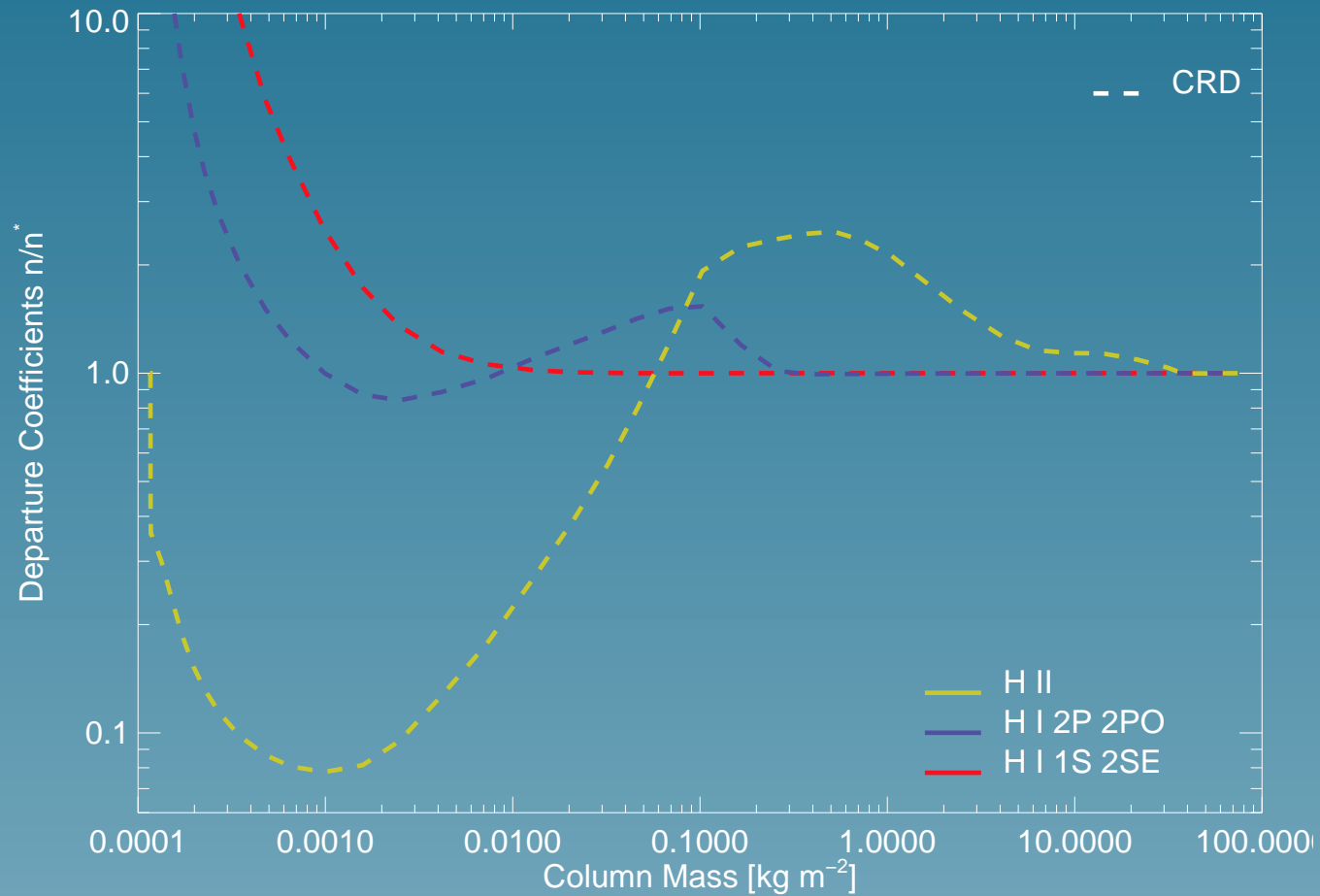
# Hydrogen Lyman- $\alpha$ profiles (quiet-Sun)



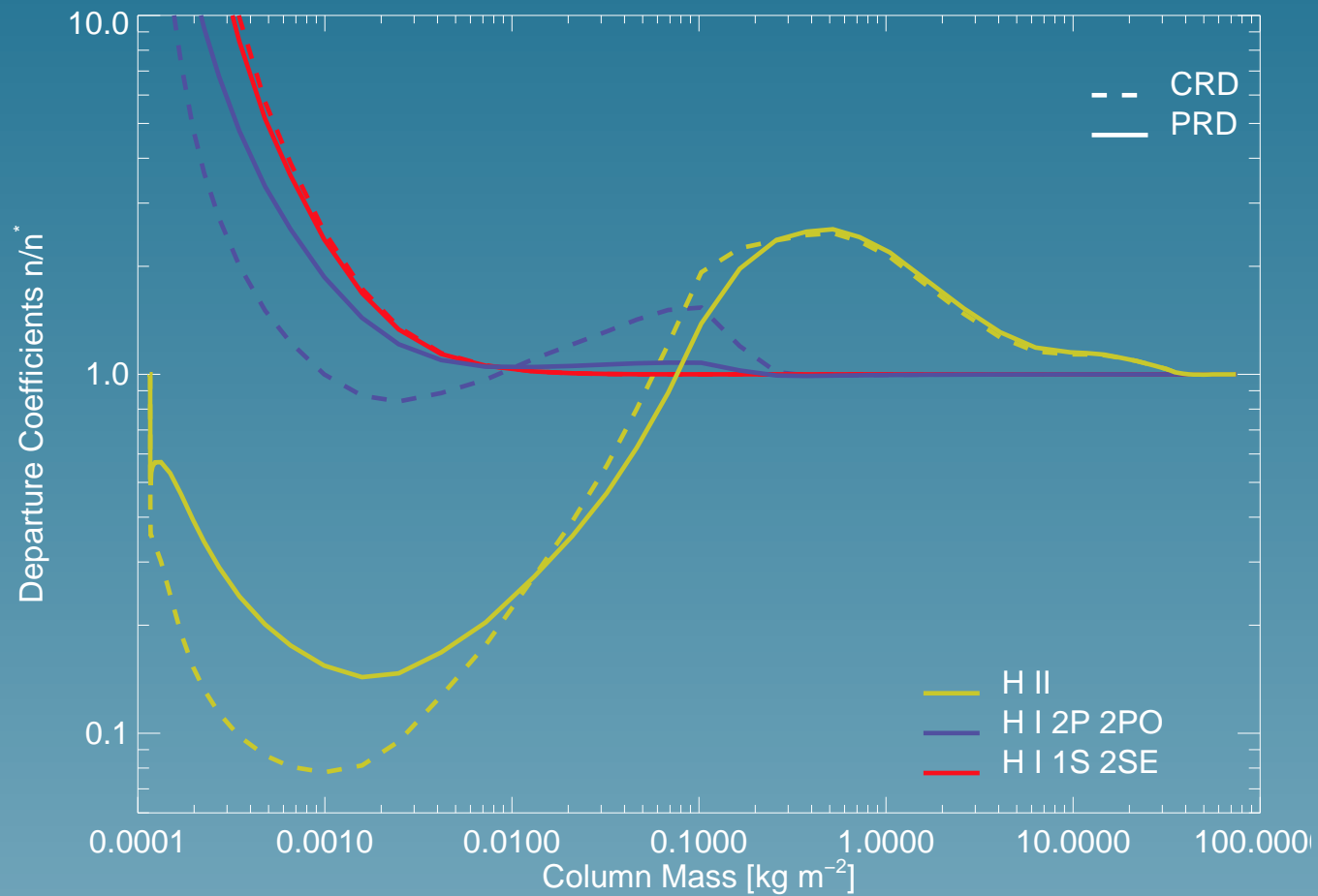
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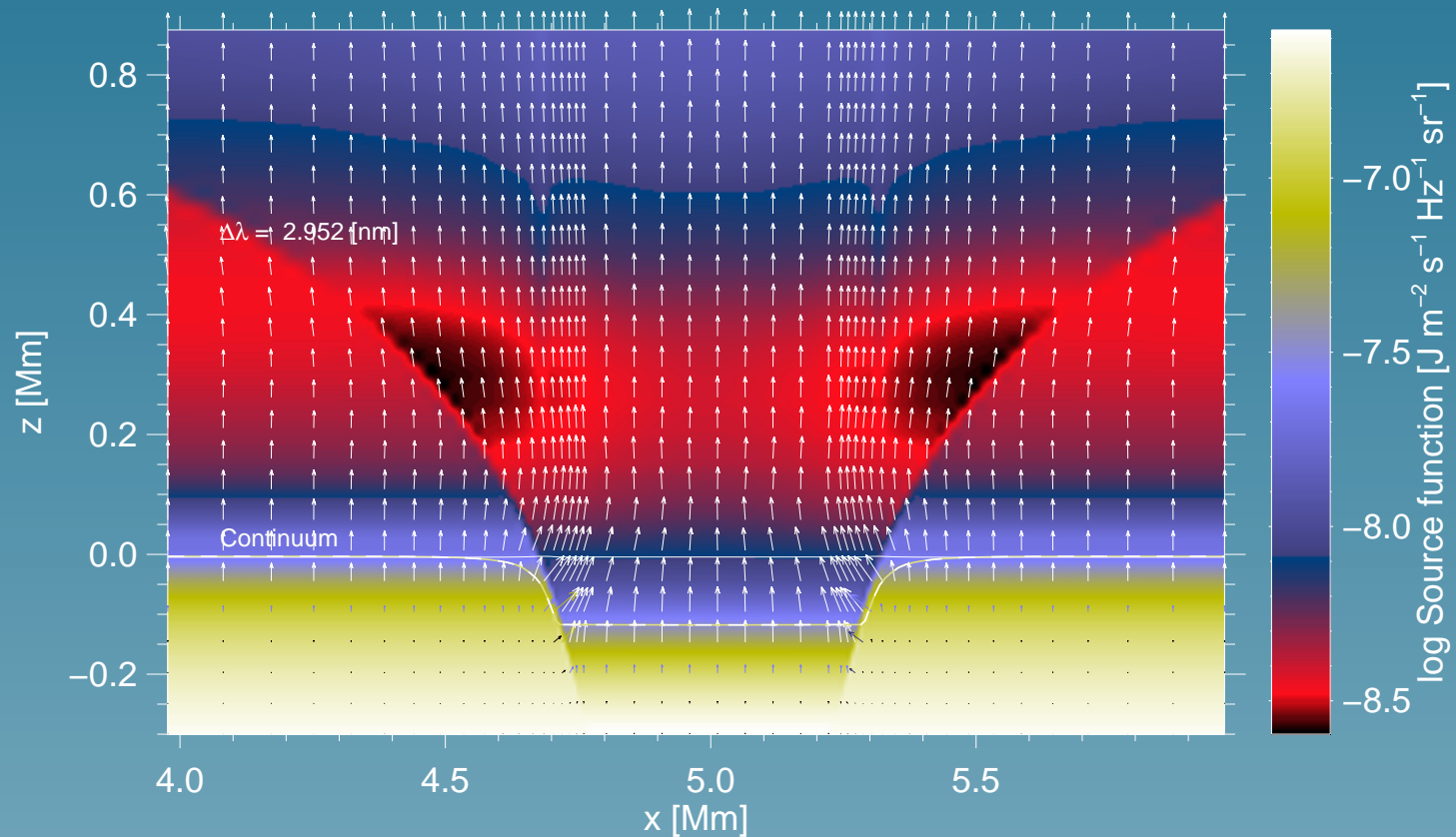
# Hydrogen Departure Coefficients



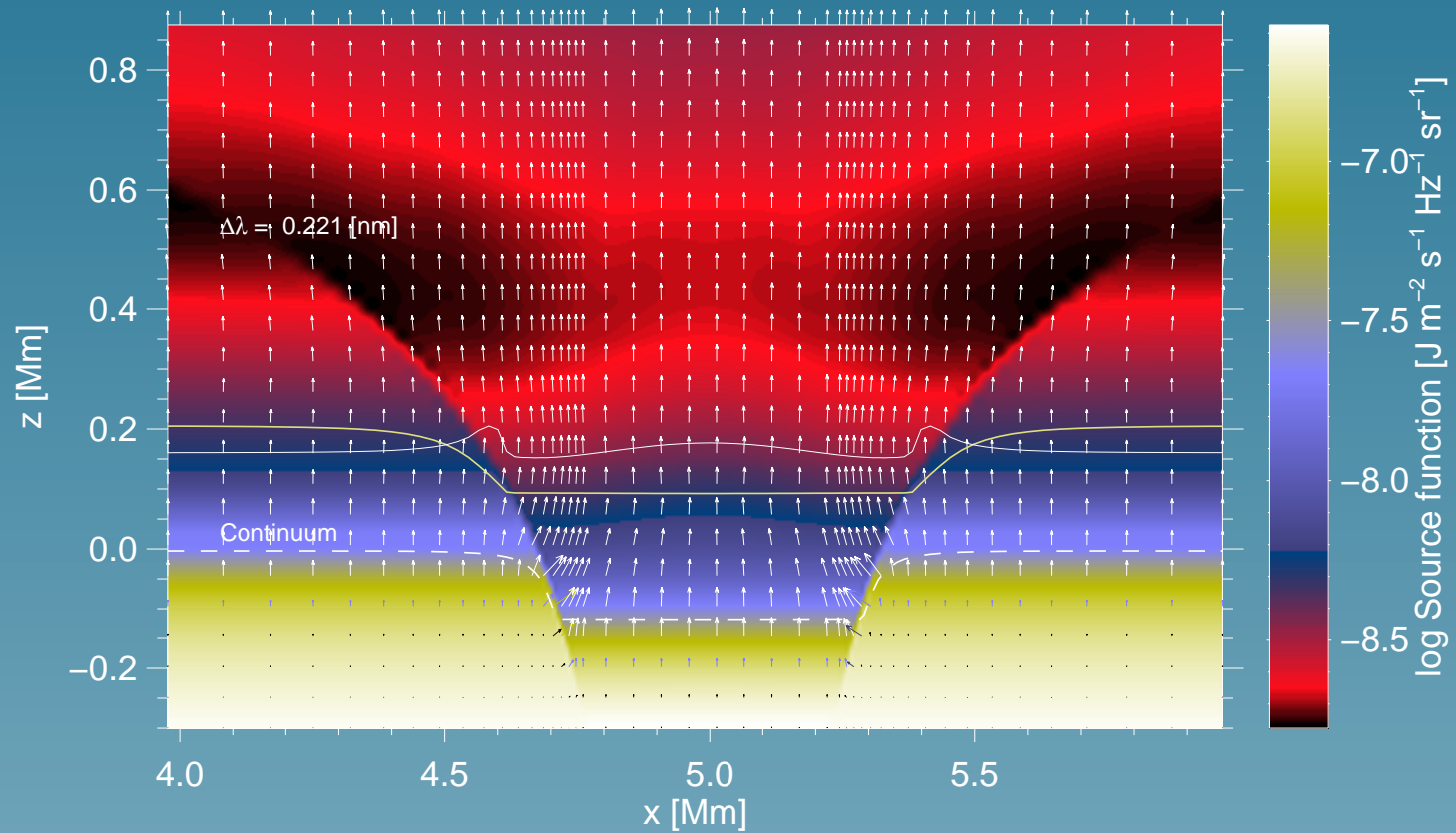
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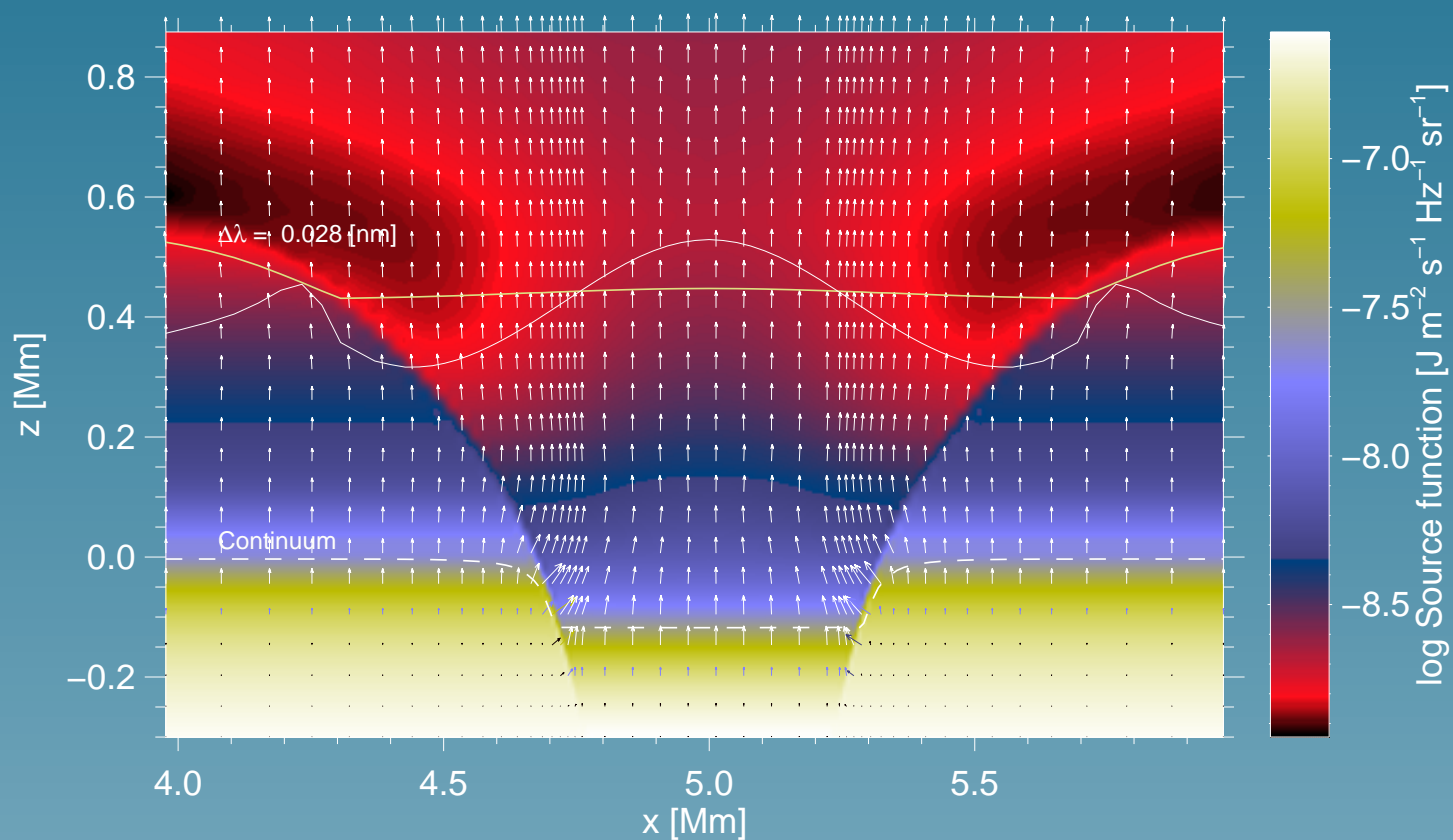
# Network Structure (Ca II)



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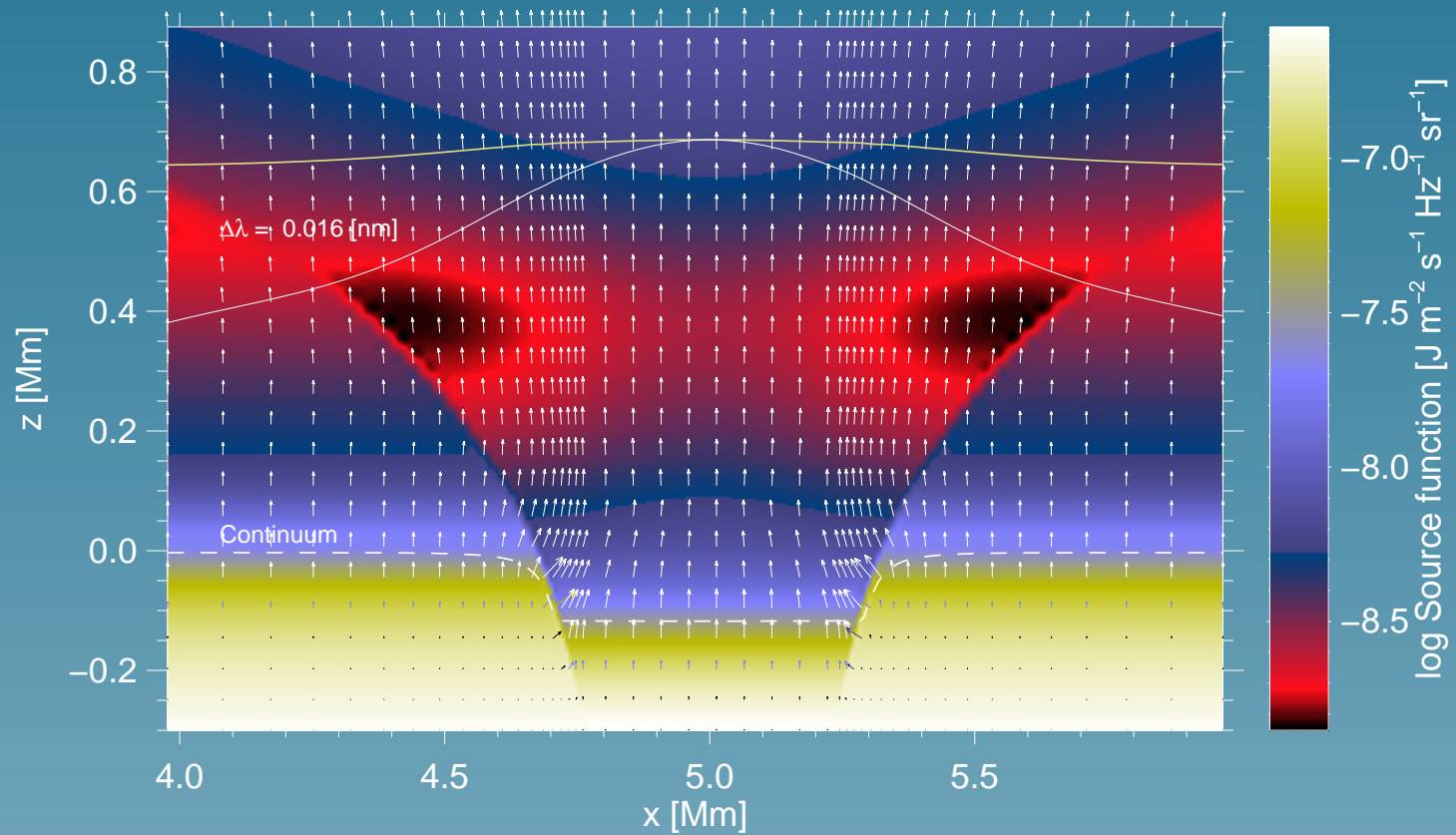


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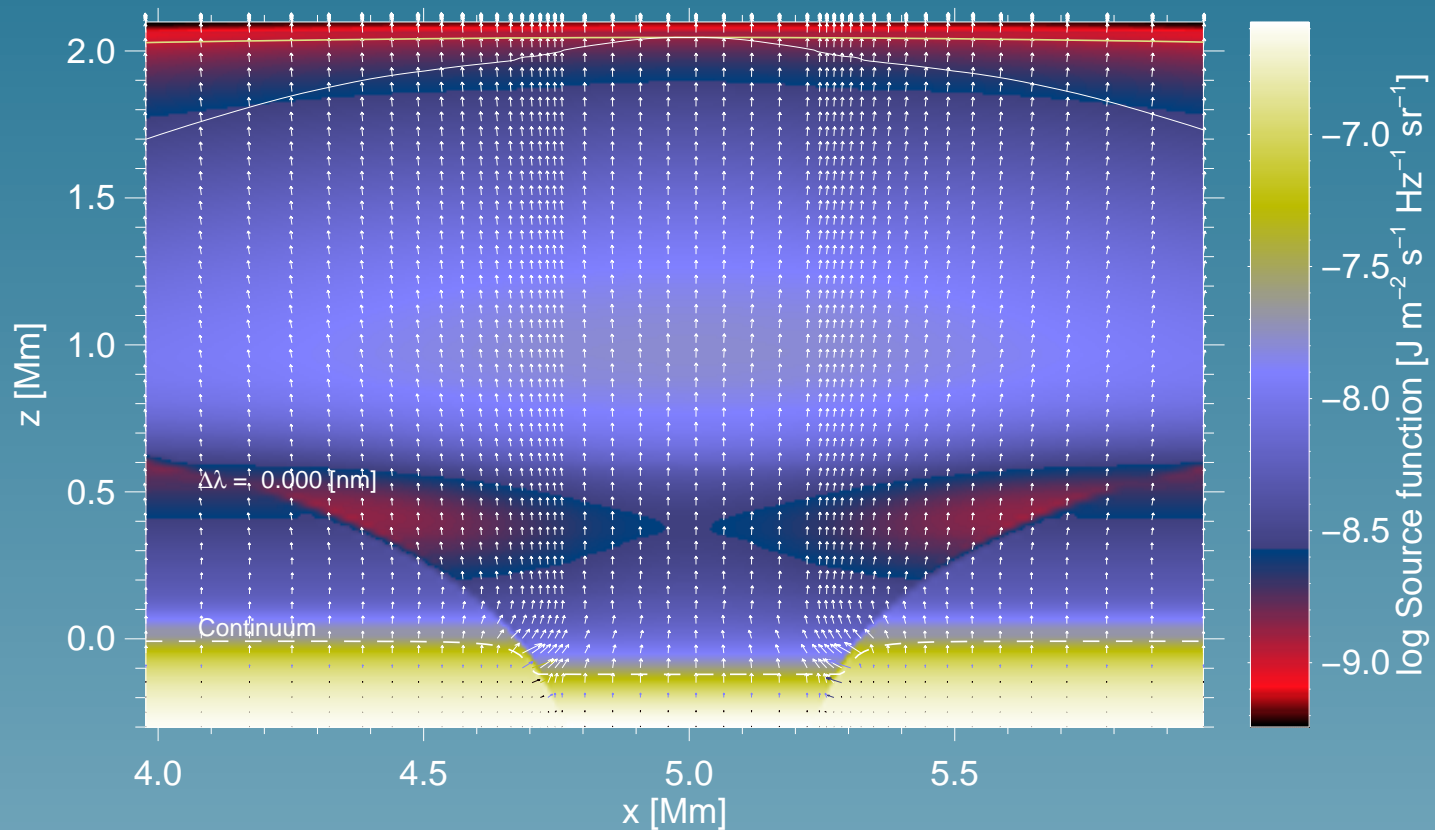




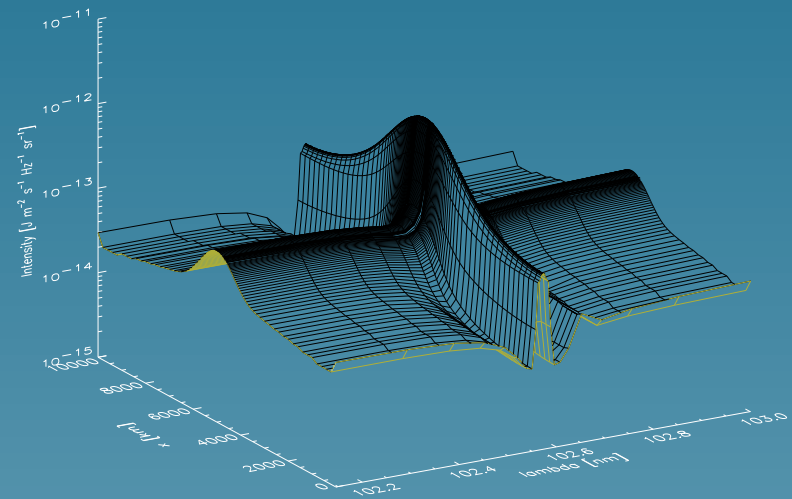
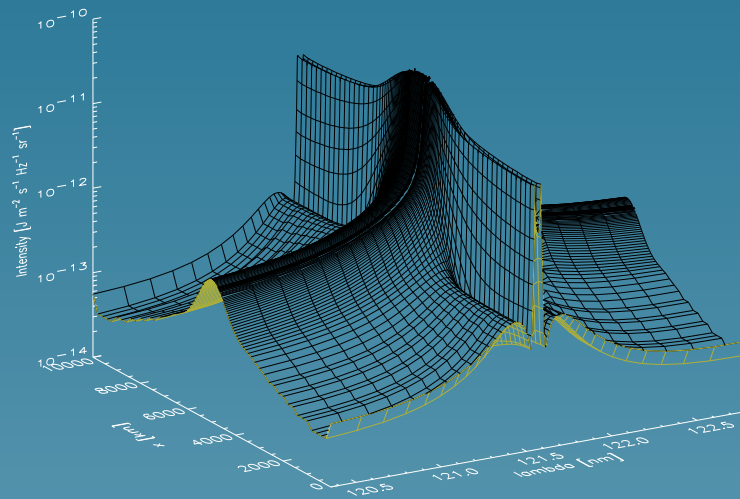
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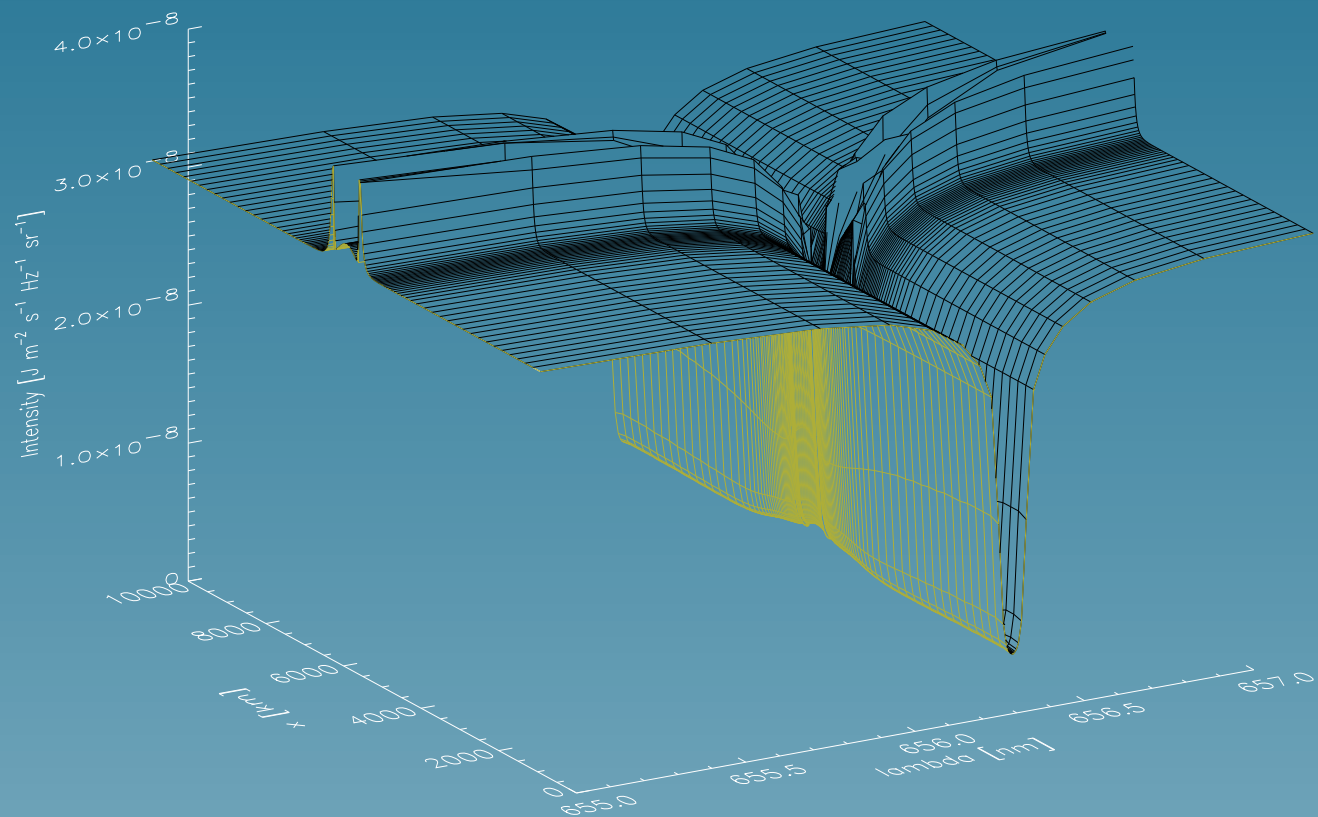
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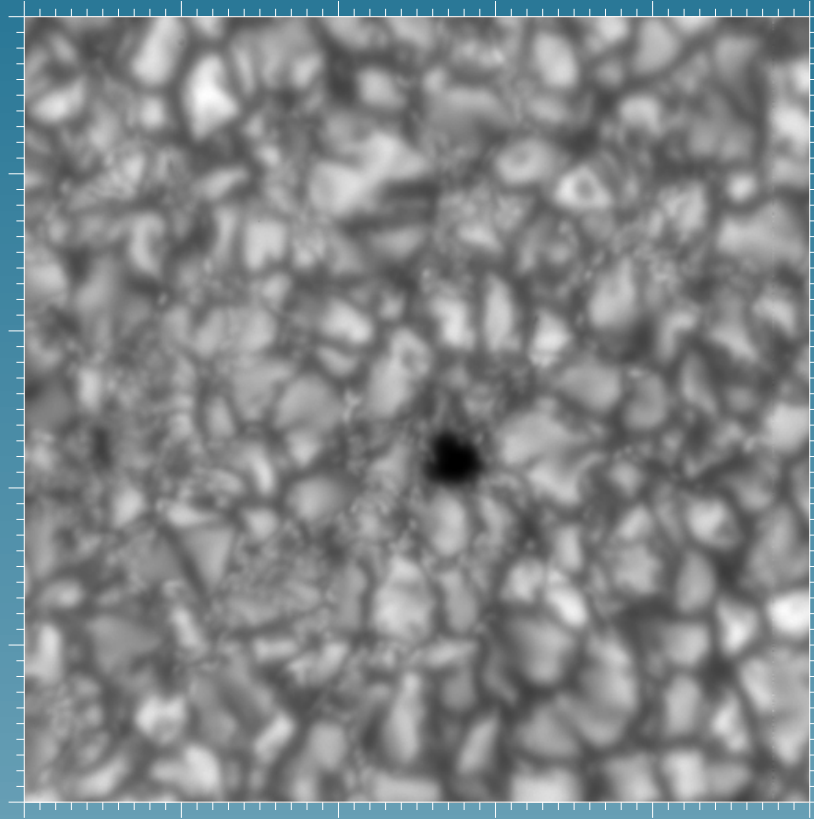
# Lyman profiles in network



# H- $\alpha$ profile in network

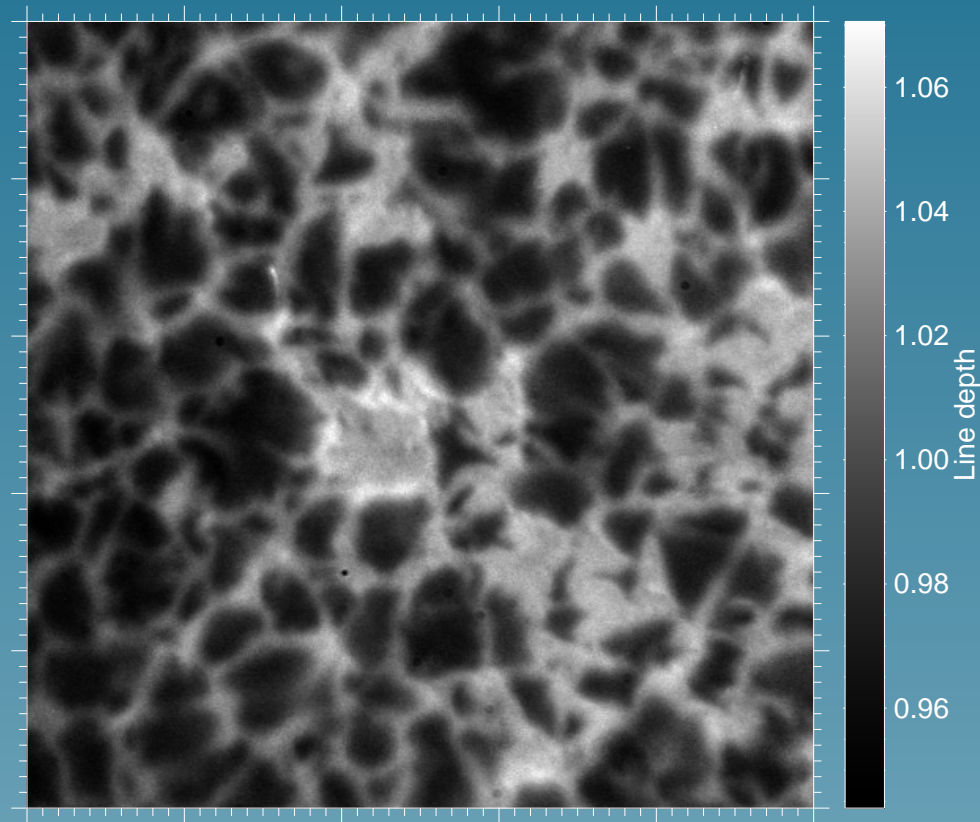


# High resolution Observations

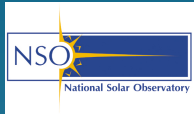


Fe I 569.1, Klaus Hartkorn and Thomas Rimmele, Sac Peak DST

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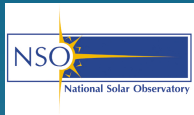


Fe I 569.1 line depth, Klaus Hartkorn and Thomas Rimmele, Sac Peak DST

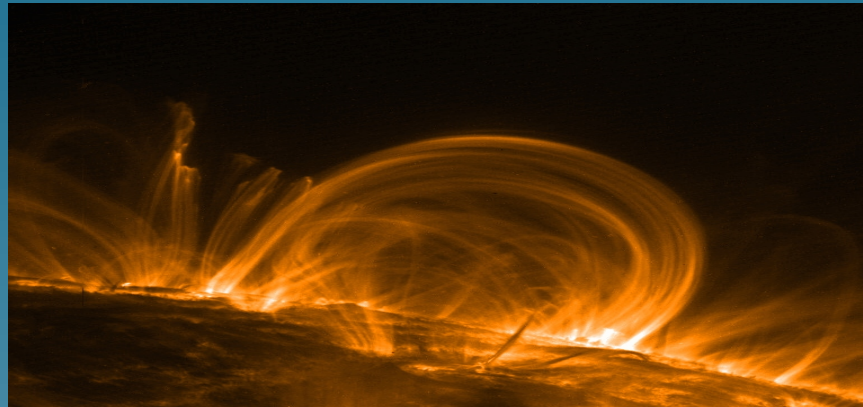


## Conclusions

- To figure out how coronal loops are heated we need to discover
  - how they are connected to the underlying photosphere
  - how energy flows from below, and in what form
- We need to develop reliable diagnostics. **Understand the information encoded in  $H\text{-}\alpha$**
- This will require multi-wavelength Stokes polarimetry with high cadence and appropriate spatial and spectral resolution



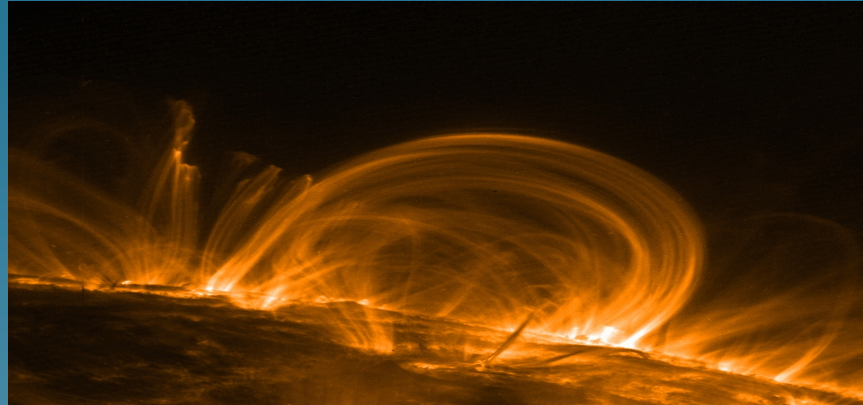
# Where are Coronal Loops Heated?



[back](#)



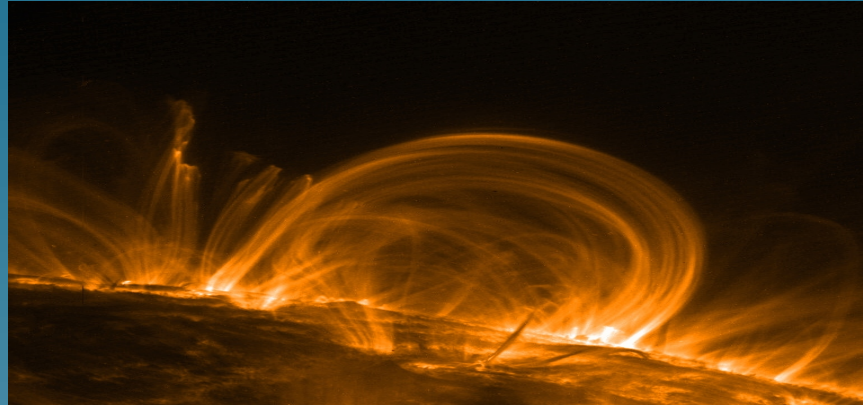
# Where are Coronal Loops Heated?



"This heating function provides uniform heating for small loops ( $L \leq 20$  Mm), but restricts heating to the footpoints of large loops ( $L \equiv 50 - 300$  Mm)." (Aschwanden et al., 2000, ApJ 541, 1059)

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## Where are Coronal Loops Heated?

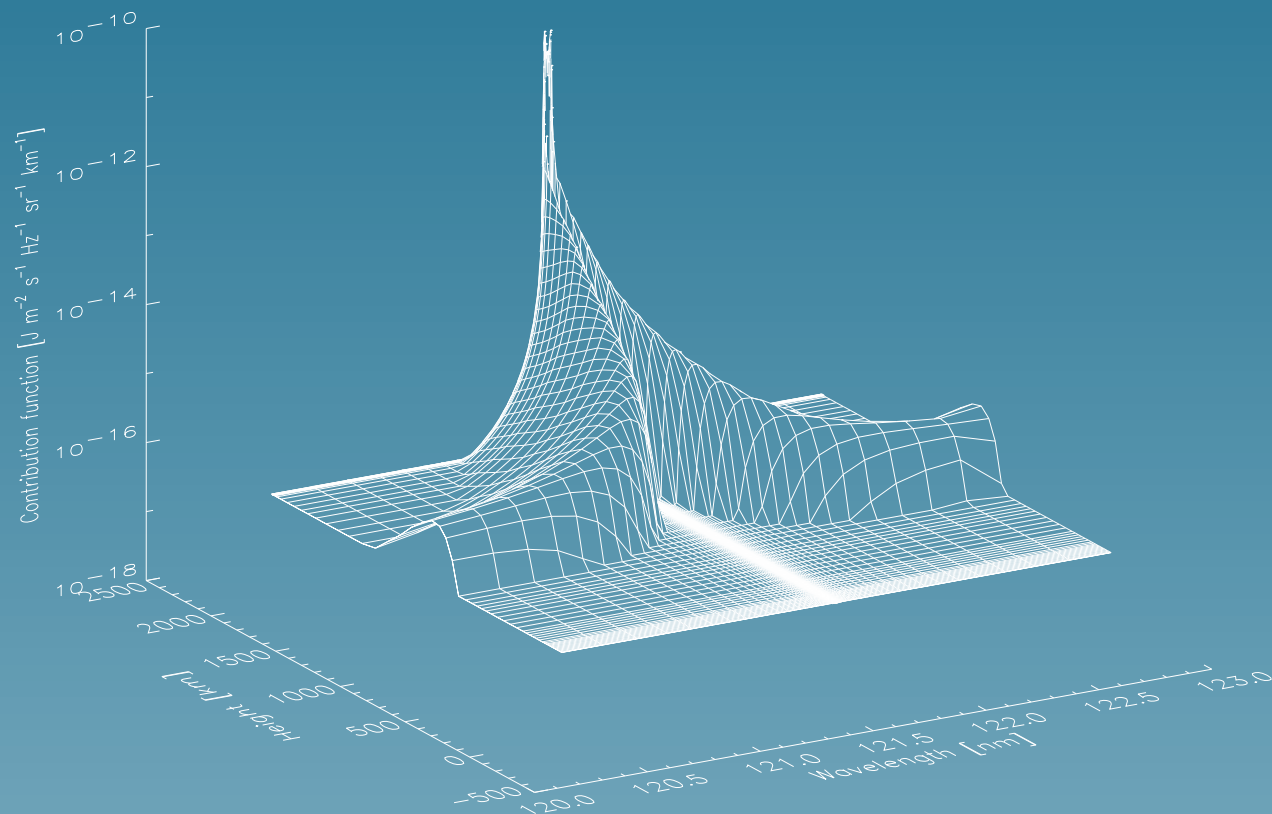


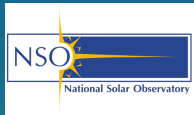
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"This gives strong evidence against heating concentrated near the loop base for such loops . . ." (Priest et al., 2000, ApJ 539, 1002)

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# Hydrogen Lyman- $\alpha$ Contribution Function





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